

CONSERVATION OF MOMENTUM

RECAP :

$$\frac{D}{Dt} \int_{\Omega} b \, d\Omega = \int_{\Omega} q^B \, d\Omega + \int_{\partial\Omega} T_{\dots j}^B n_j \, dA \quad \dots \text{integral form of conservation law for } B.$$

This video : $B =$ momentum (rank 1 tensor i.e. a vector).

$$\begin{aligned} b &= \rho \underline{u} && \text{(density of momentum)} \\ q^B &= \rho \underline{g} && \text{volumetric force density, a.k.a body force} \end{aligned} \quad \left| \begin{array}{l} \text{Rate of change of momentum} \\ \text{is a force, so the generation} \\ \text{rates are all forces.} \end{array} \right.$$

e.g. the force of gravity, where $\underline{g} =$ acceleration due to gravity.

$T_{ij} = T_{ji}^B =$ Cauchy stress tensor, a second rank tensor that can be used to represent forces transmitted across surfaces.

Conservation of momentum in integral form

$$\frac{D}{Dt} \int_{\Omega} \rho \underline{u} \, d\Omega = \int_{\Omega} \rho \underline{g} \, d\Omega + \int_{\partial\Omega} \underline{T} \cdot \underline{\hat{n}} \, dA \quad \text{OR} \quad \frac{D}{Dt} \int_{\Omega} \rho u_i \, d\Omega = \int_{\Omega} \rho g_i \, d\Omega + \int_{\partial\Omega} T_{ij} n_j \, dA.$$

Using Reynolds transport theorem for the L.H.S.

$$\begin{aligned} \int_{\Omega} \frac{\partial(\rho \underline{u})}{\partial t} \, d\Omega + \int_{\partial\Omega} \rho \underline{u} (\underline{u} \cdot \underline{\hat{n}}) \, dA &= \int_{\Omega} \rho \underline{g} \, d\Omega + \int_{\partial\Omega} \underline{T} \cdot \underline{\hat{n}} \, dA \\ \text{OR} \quad \int_{\Omega} \frac{\partial(\rho u_i)}{\partial t} \, d\Omega + \int_{\partial\Omega} \rho u_i (u_j n_j) \, dA &= \int_{\Omega} \rho g_i \, d\Omega + \int_{\partial\Omega} T_{ij} n_j \, dA. \end{aligned}$$

Using divergence theorem

$$\int_{\Omega} \left[\frac{\partial(\rho \underline{u})}{\partial t} + \nabla \cdot (\rho \underline{u} \underline{u}) - \rho \underline{g} - \nabla \cdot \underline{T} \right] d\Omega = 0 \quad \text{OR} \quad \int_{\Omega} \left[\frac{\partial(\rho u_i)}{\partial t} + \frac{\partial(\rho u_i u_j)}{\partial x_j} - \rho g_i - \frac{\partial T_{ij}}{\partial x_j} \right] d\Omega = 0$$

Conservation of momentum in differential form

$$\frac{\partial}{\partial t}(\rho \underline{u}) + \nabla \cdot (\rho \underline{u} \underline{u}) = \rho \underline{g} + \nabla \cdot \underline{T} \quad \text{OR} \quad \frac{\partial}{\partial t}(\rho u_i) + \frac{\partial}{\partial x_j}(\rho u_i u_j) = \rho g_i - \frac{\partial T_{ij}}{\partial x_j}.$$

Note that $\frac{\partial}{\partial t}(\rho u_i) + \frac{\partial}{\partial x_j}(\rho u_i u_j) = \underbrace{\rho \left[\frac{\partial u_i}{\partial t} + u_j \frac{\partial u_i}{\partial x_j} \right]}_{\text{acceleration } \underline{a}} + u_i \left[\frac{\partial \rho}{\partial t} + \frac{\partial(\rho u_j)}{\partial x_j} \right]$ ↑ mass conservation

$$\Rightarrow \rho \frac{D \underline{u}}{Dt} = \rho \underline{g} + \underbrace{\nabla \cdot \underline{T}}_{\substack{\text{surface force} \\ \text{per unit volume} \\ \text{on an infinitesimal volume.}}} \quad \text{OR} \quad \rho \frac{D u_i}{Dt} = \rho g_i + \frac{\partial T_{ij}}{\partial x_j}$$

